

$$\mu = \frac{D}{k_B T} = \frac{1}{6\pi\eta \cdot a}$$

Langevin-Equation

- Stokes law = $F_S = -\gamma \cdot \dot{x}$
- Newton's 2nd law $F_N = m \ddot{x}$

$$m \ddot{x} = -\gamma \dot{x} + F_{\text{ext}} + \gamma(t)$$

↑ noise

$$F_{\text{ext}} = 0$$

$$V(x) = V_0 \exp\left(-\frac{\zeta}{\lambda} \cdot x\right)$$

↑ damping term

- zero mean $\langle \gamma(t) \rangle = 0$
- gaussian distributed
- white noise

$$\langle \gamma(t) \cdot \gamma(t') \rangle = 2\gamma k_B T \delta(t - t')$$

overdamped Langevin Eq.

$$\ddot{x} = 0$$

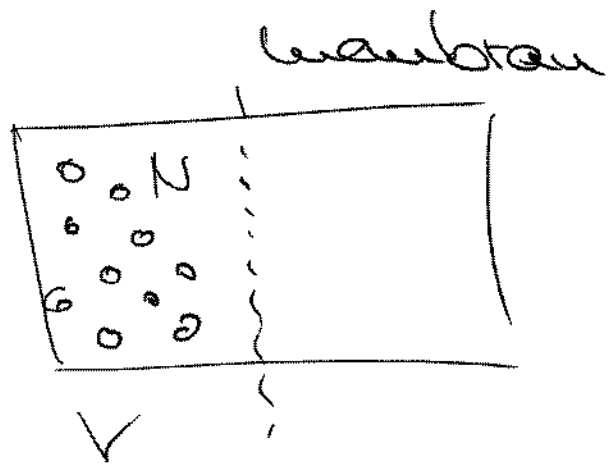
$$\gamma \dot{x} = F_{\text{ext}} + \gamma(t)$$

Stokes - Einstein - Relation

$$\vec{f} = \vec{f}_D + \vec{f}_{\text{fied}} \stackrel{!}{=} 0 \quad \text{equilib.}$$

$$-\eta \nabla^2 \vec{s}_{\text{eq}} + \mu \vec{F} \cdot \vec{s}_{\text{eq}} = 0$$

$\vec{s}_{\text{eq}}?$



$$\rho = \frac{N}{V}$$

in equilib.

$$P(r) = \rho(r) \cdot k_B T$$

$$F = m \cdot g$$

$$\rho \cdot F + \nabla^2 P(r) = 0$$

$$\frac{P(r) \cdot F(r)}{k_B T} = -\nabla^2 P(r) \rightarrow \rho_{\text{eq}} \propto P(r) \propto \exp\left(\frac{-F \cdot r}{k_B T}\right)$$

$$\vec{J}_{\text{field}} = \vec{v} \cdot g$$

↑
mean drift velocity

$$\vec{v} = \mu \vec{E}$$

↑
mobility

$$m \dot{\vec{v}} + \gamma \vec{v} = \vec{F}$$

↑
friction coefficient

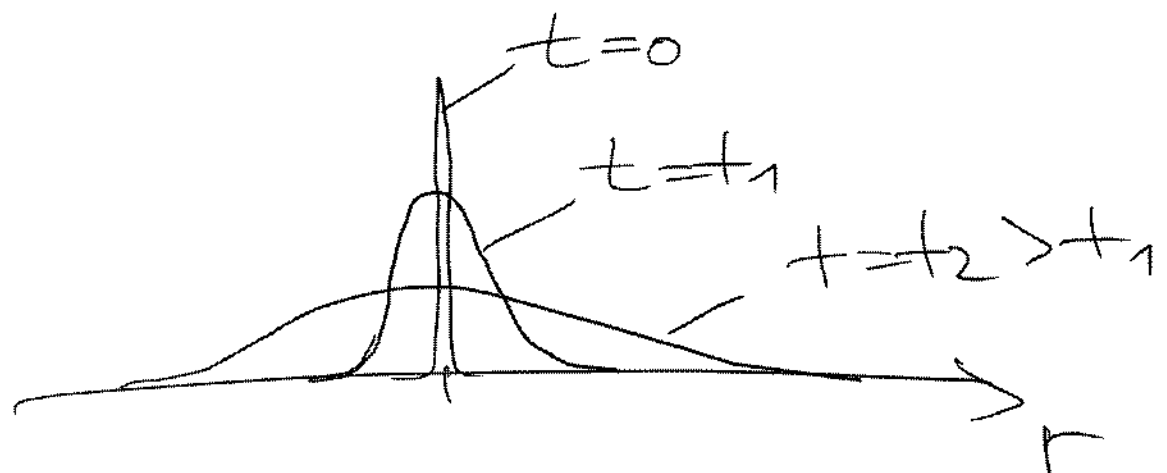
Stationary case $\dot{\vec{v}} = 0$

$$\vec{v} = \vec{F} / \gamma$$

$$\rightarrow \vec{v} = \mu \cdot \vec{F} \cdot g$$

$$\mu = \frac{1}{\gamma}$$

$$\vec{F} = \frac{-6\pi\eta a v}{\gamma}$$



broadering $\hat{=}$ dispersion

mean-square-displacement (MSD)

$$\langle (r(t) - r(0))^2 \rangle = \int d^3r g(\vec{r}, t) (r(t) - r(0))^2$$

$$= 6Dt$$

now apply external force



diffusion + drift (lin. superposition)

continuity equation

$$\frac{\partial g(\vec{r}, t)}{\partial t} = -\vec{\nabla} \cdot \vec{j}(\vec{r}, t)$$

→ diffusion equation

$$\frac{\partial g(\vec{r}, t)}{\partial t} = D \Delta g(\vec{r}, t)$$

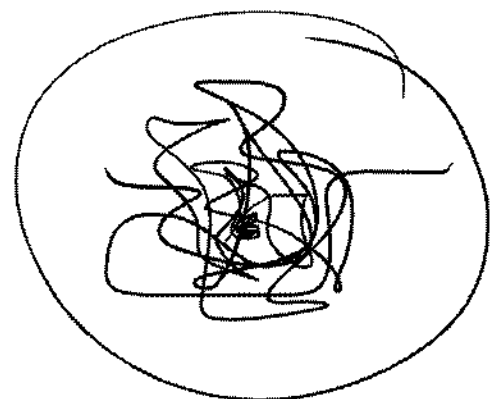
Starting condition $g(\vec{r}, 0) = \delta(\vec{r})$

$$g(\vec{r}, t) = (4\pi Dt)^{-3/2} \cdot e^{-\frac{(\vec{r}(t) - \vec{r}(0))^2}{4Dt}}$$

gaussian
distribution

Historical overview

Solution \rightarrow A. Einstein



trajectoire

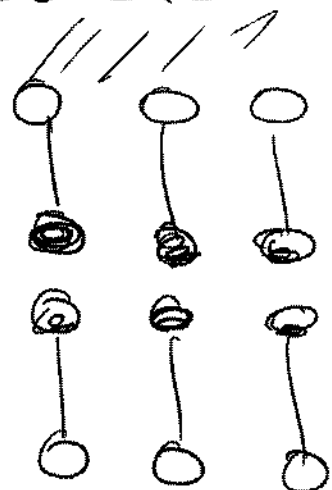
random walk

Langevin - theory of Brownian motion

$$\vec{f}_D(\vec{r}, t) = -D \vec{\nabla} g(\vec{r}, t)$$

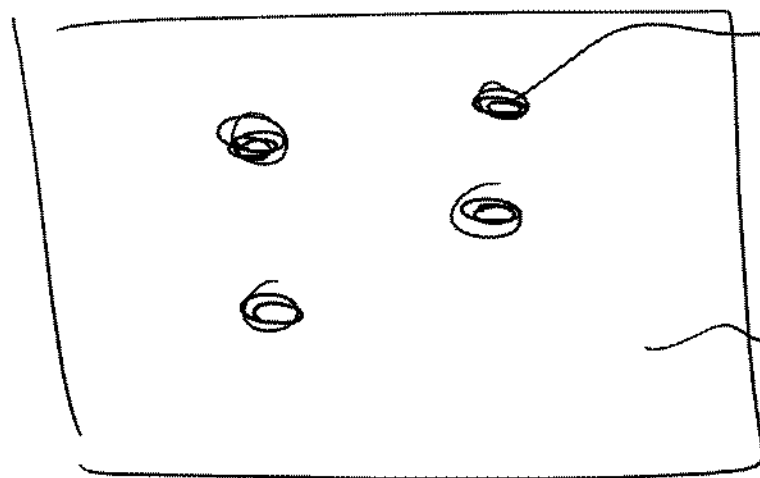
↑
diffusion coefficient

5) vesicles



≙ biolog. cell membrane

→ back to colloids

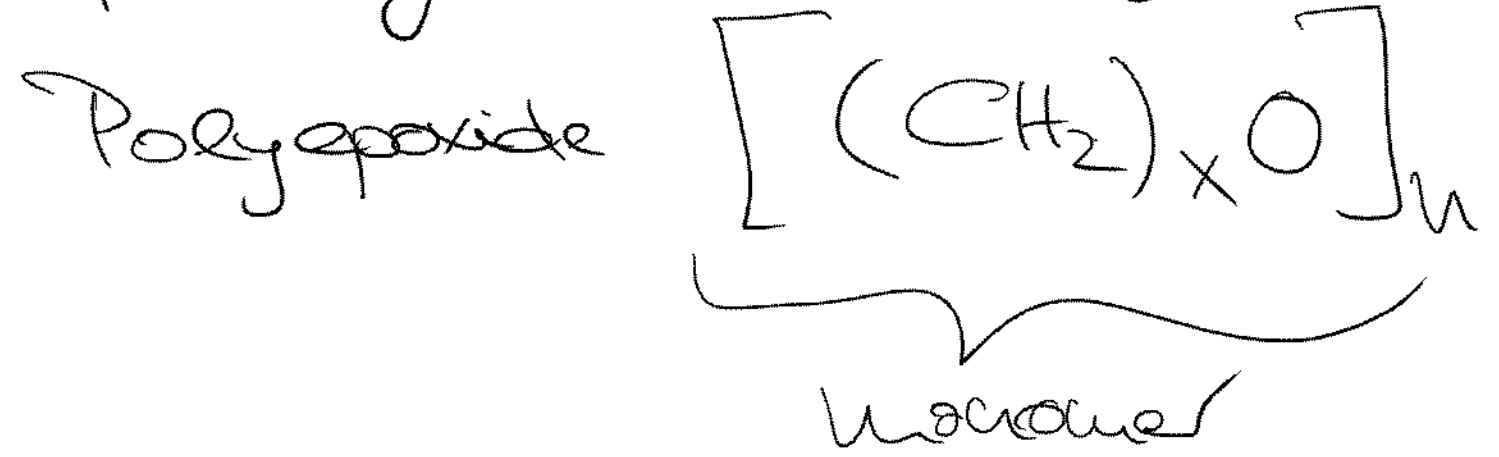


dispersed phase

solvent phase

3) Polymers

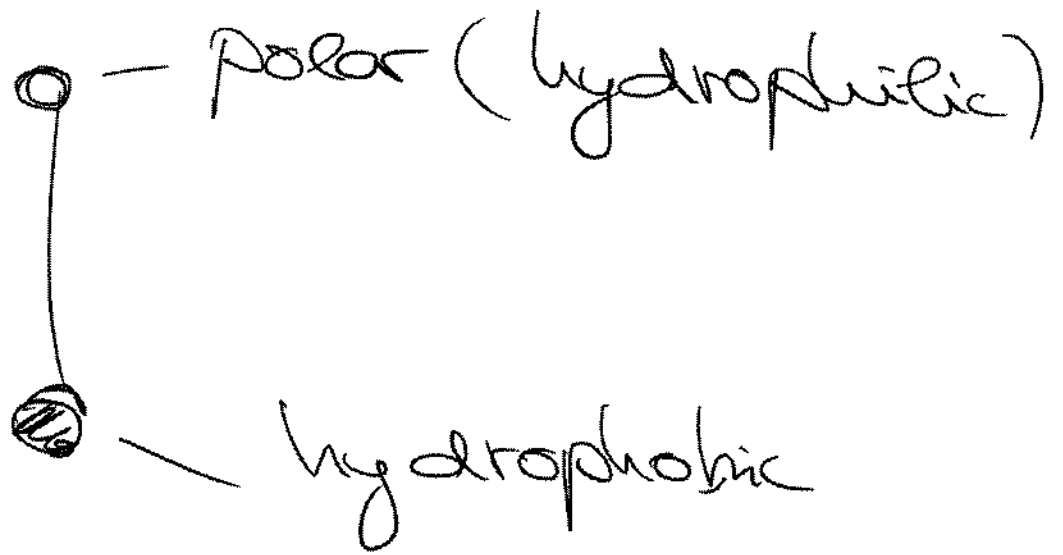
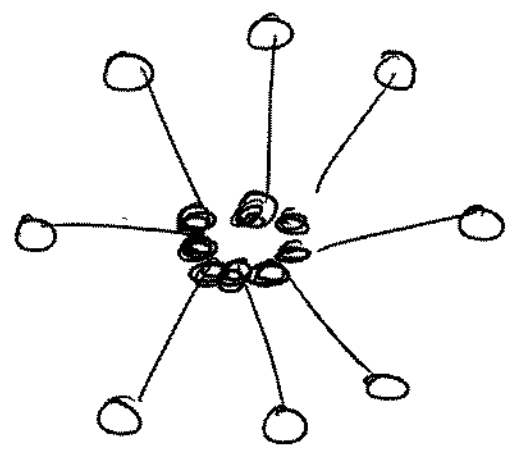
repeating unit consisting of 100 - 10.000 atoms



4) micelles

amphiphilic molecule

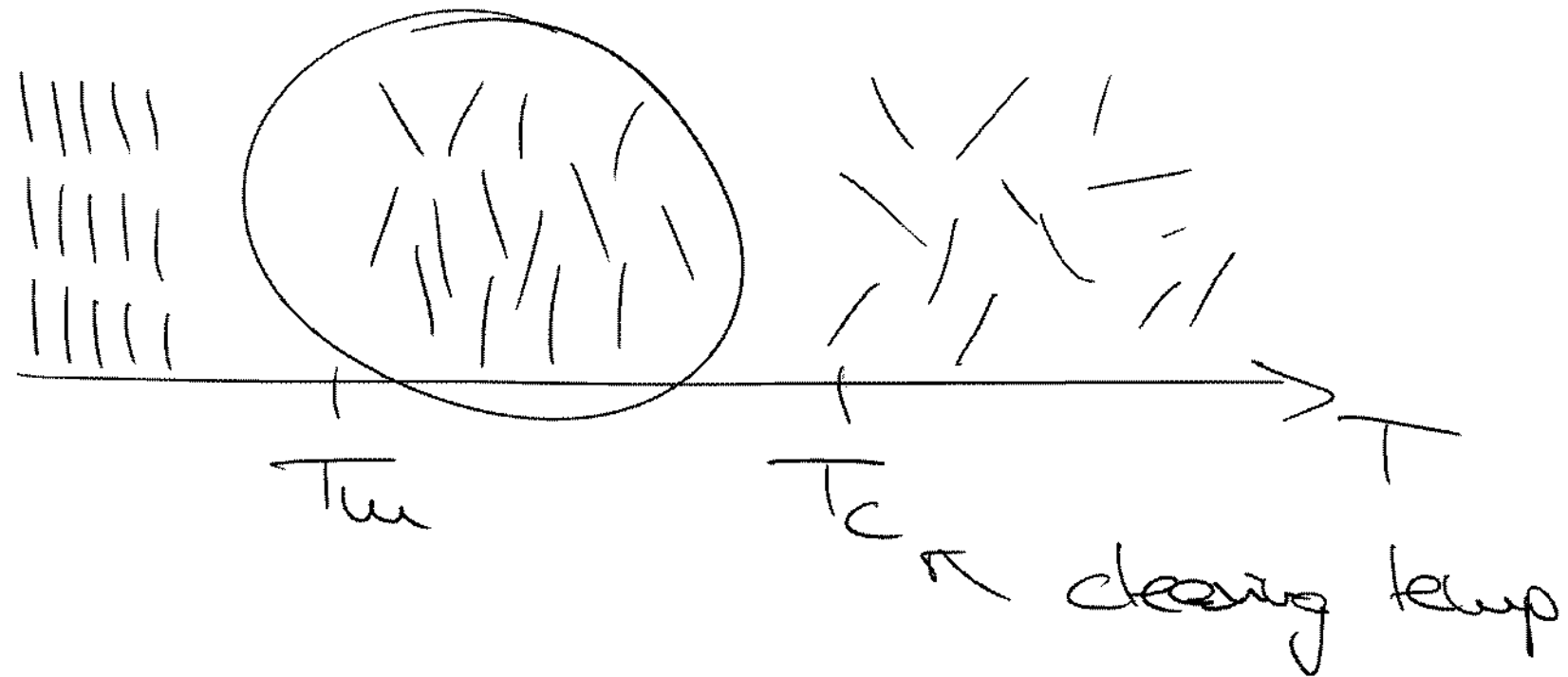
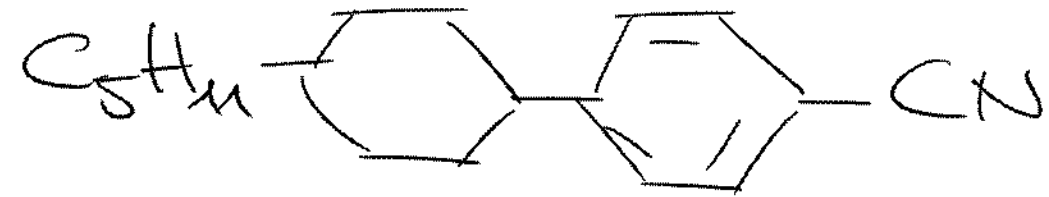
polar solvent



2) Liquid crystals

anisotropic molecules $\sim 10 - 100$ atoms

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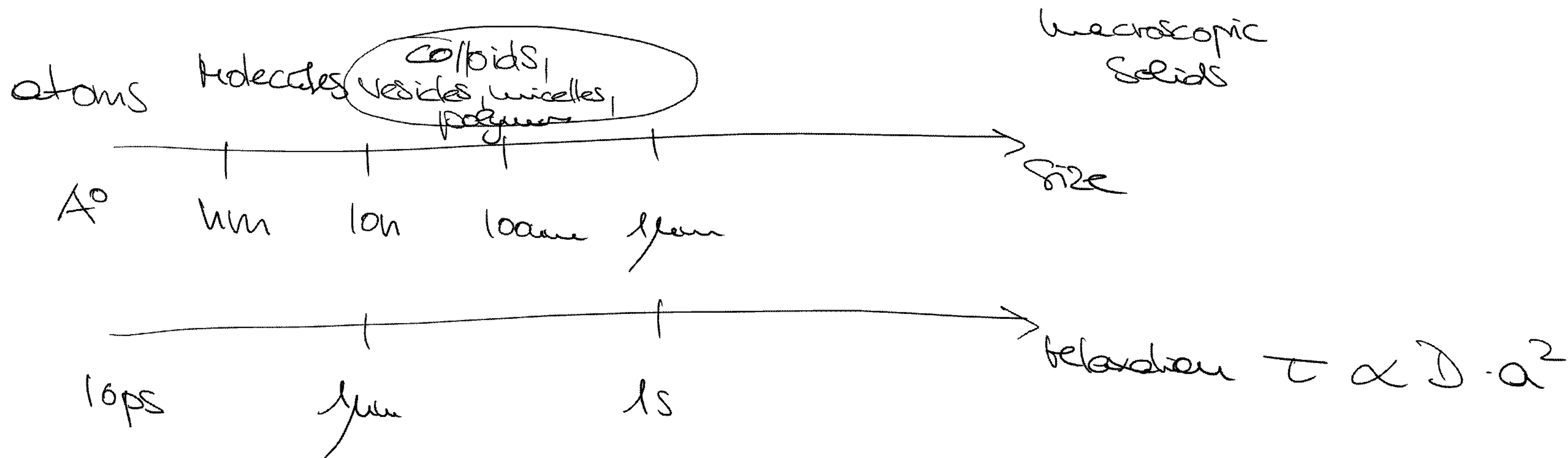
elastic constant: $\frac{k_B T}{a_{\text{Koll}}^3} / \frac{k_B T}{a_{\text{Atom}}^3} \ll 10^{-9}$

→ macromolecules are easily deformable
high compressibility

1) colloidal systems

Solid particles suspended in liquid
 $a \sim 10 \text{ nm} - 10 \mu\text{m}$

What is soft matter?



- $\frac{\Omega_{\text{coll}}}{\Omega_{\text{atom}}} \geq 10^3$; typ. interaction energy between particles $\sim k_B T \sim \frac{1}{40} \text{ eV}$