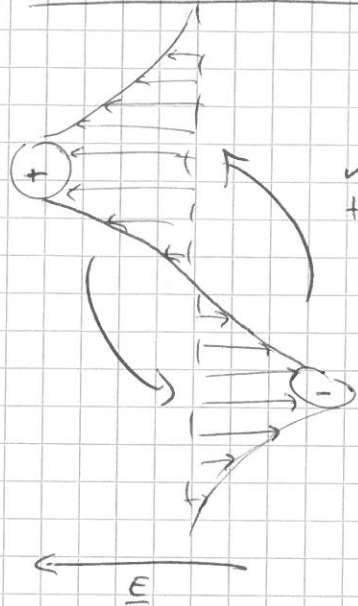
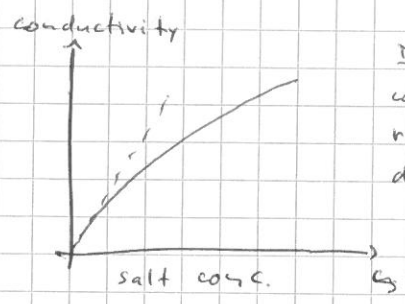


Hydrodynamic interactions

Video sedimentation instability



momentum exchange through fluid



Decrease in cond. of electrolyte rel. to c_s is partially due to hyd. int.

$Sc = \frac{\tau_{kin}}{D}$ "ratio of momentum transport to mass transport"

$\approx 10^{-2} - 10^{-3}$ for molecular liquids

$\approx 10^6$ for μm -sized colloids

Momentum transport is typically much faster than mass transport (particles). Flow can be assumed stationary on small (μm) scales.

Governing equations for fluid flow

Momentum is a conserved quantity. Special relativity dictates that there must be a continuity eq.

$$\nabla \cdot \underline{\underline{\Pi}} = -\partial_t (\underline{\underline{S}} \underline{\underline{u}}) + \underline{\underline{f}}$$

stress tensor density velocity external force

momentum flux momentum density $\underline{\underline{\Pi}} \underline{\underline{u}}$ "momentum transported in direction $\underline{\underline{u}}$ "

momentum sources

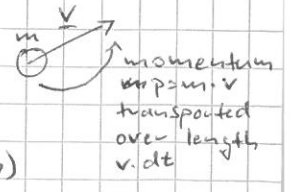
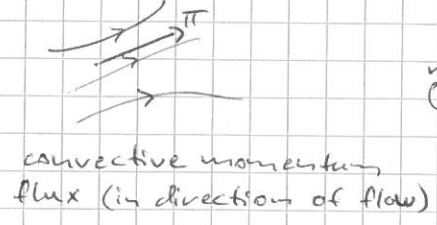
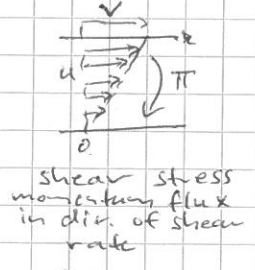
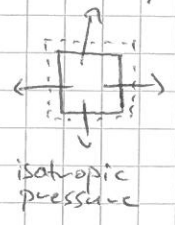
c.f. $\nabla \cdot \underline{\underline{j}} = -\partial_t c + R$

flux density source/sinks $\underline{\underline{j}} \cdot \underline{\underline{u}}$ flux in dir. $\underline{\underline{u}}$

For an incompressible, Newtonian fluid

$$\underline{\underline{\Pi}} = p \underline{\underline{1}} - \eta (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) + \underline{\underline{S}} \underline{\underline{u}} \otimes \underline{\underline{u}}$$

$\underline{\underline{u}} \otimes \underline{\underline{u}}$ is projector onto $\underline{\underline{u}}$ and quadratic in $\underline{\underline{u}}$



$\eta_{air} < \eta_{water} < \eta_{oil} < \eta_{tar}$

Plugging Π into continuity eq., using incompressibility $\nabla \cdot \underline{u} = 0$, $S = \text{const}$, yields (after some tensor-calculus)

$$\boxed{\begin{aligned} S \left[\partial_t \underline{u} + (\underline{u} \cdot \nabla) \underline{u} \right] &= -\nabla p + \eta \nabla^2 \underline{u} + \underline{f} \\ \nabla \cdot \underline{u} &= 0 \end{aligned}}$$

incompressible Navier-Stokes Eq. for a Newtonian fluid

De-dimensionalization with characteristic length a , velocity v and time τ yields

$$\frac{S a^2}{\eta \tau} \partial_t \underline{u}' + \text{Re} (\underline{u}' \cdot \nabla) \underline{u}' = -\nabla p' + \nabla^2 \underline{u}' + \underline{f}'$$

$$\text{Re} = \frac{S v a}{\eta} \quad \text{"ratio of convective to viscous momentum transport"}$$

$\approx 10^5$ for macroscopic systems (water at 20°C, 0.1 m sphere falling at 1 m/s)

$\approx 10^{-5}$ for colloidal system or bacteria ($a = 1 \mu\text{m}$, $v = 10 \mu\text{m/s}$)

Similarity principle: Flows in the same geometry and with the same Re are similar.

Different flow videos

"Small animals experience swimming in water like humans experience swimming in honey or even tar."

In typical soft-matter systems $\text{Re} \ll 1$, which allows us to use the simpler Stokes Eq.

$$\boxed{\begin{aligned} \eta \nabla^2 \underline{u} &= \nabla p - \underline{f} \\ \nabla \cdot \underline{u} &= 0 \end{aligned}}$$

Stokes Eq.

Since this eq. is linear, it has a useful Green's function

$$\underline{f} = \underline{F} \cdot \delta(\underline{r}) \Rightarrow \begin{cases} \underline{u}(\underline{r}) = \frac{\underline{F}}{8\pi\eta} \left(\frac{1}{r} \underline{1} + \frac{1}{r^3} \underline{r} \otimes \underline{r} \right) \\ p(\underline{r}) = \frac{1}{4\pi r^3} \underline{F} \cdot \underline{r} \end{cases} \quad \text{Oseen-Tensor}$$

There are improvements over the Oseen-Tensor, e.g. the Rotney-Prager-Tensor, which assumes spheres with no-slip BC and incorporates rotations.

This allows one to treat HI as pair-forces. These tensors yield a linear relation between all particle forces and all particle velocities. One can determine the velocities given all conservative forces. ~~by solving a dense linear system of eq. in every time step~~

→ Stokesian dynamics, Brownian dynamics.